

Numerical Optimization of Single-Mode Photonic Crystal VCSELs

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Outline

Introduction

- ❑ Photonic crystal (PhC) VCSELs
- ❑ Single-mode condition defined by waveguiding
- ❑ Loss discrimination effect

3-D coupled simulation model

- ❑ Electro-thermal models
- ❑ Direct solution of Helmholtz equation
- ❑ Rate-equation approach

Results

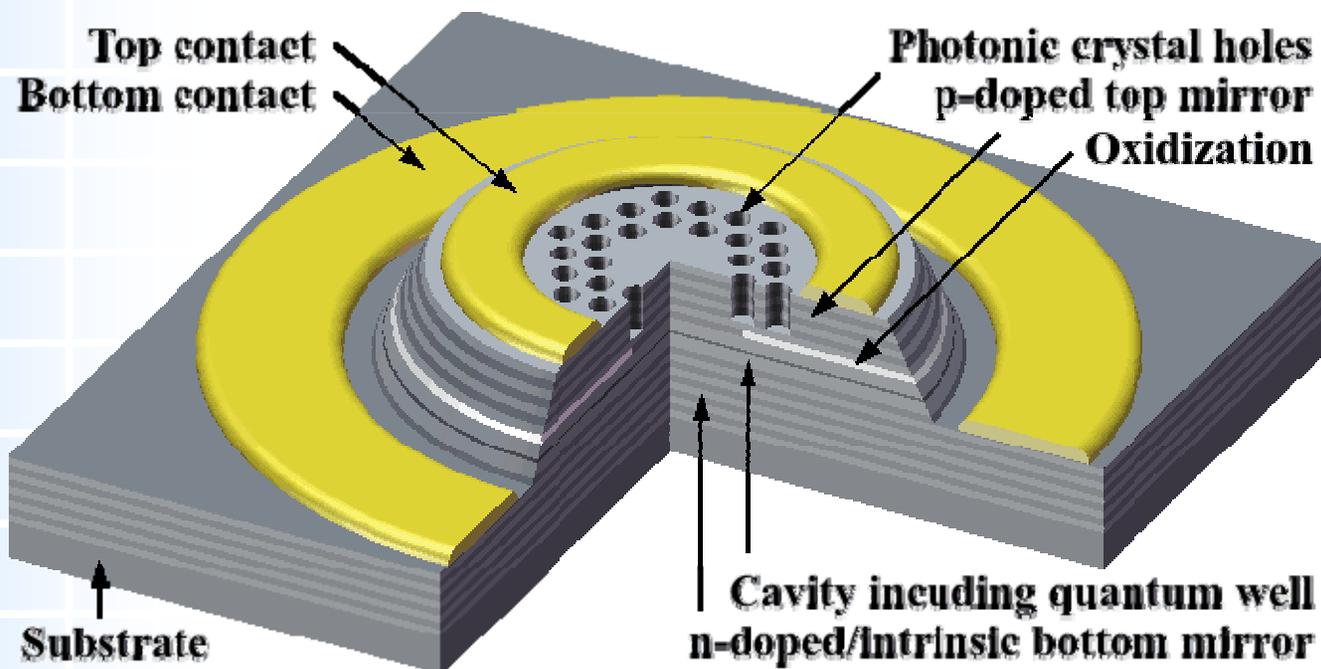
- ❑ Current and temperature distributions
- ❑ Optical mode profiles and loss discrimination
- ❑ Static light power versus current (L-I) diagrams

Summary and outlook

Photonic Crystal VCSELs

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- ❑ Optical confinement: PhC and optionally oxide aperture
- ❑ Electrical confinement: proton-implant or oxide aperture
- ❑ They potentially offer high-power single-mode operation.
- ❑ Design parameters: lattice constant (Λ , a), hole diameter (d) or diameter-to-pitch ratio, etching depth, and the diameter of the electrical aperture



Single-Mode PhC-VCSELS

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Normalized frequency parameter

$$V_{\text{eff}} = 2\pi r / \lambda \sqrt{n_{\text{eff}}^2 - (n_{\text{eff}} - \gamma \Delta n)^2} \quad [1]$$

- n_{eff} : refractivity of core
- Δn : index change introduced by full PhC
- γ : etching depth factor (0-1)
- $V_{\text{eff}} < 2.405$ corresponds to the single-mode regime [2]

Normalized propagation constant

- 2-D Helmholtz equation for the transverse cross section (fully etched VCSEL)

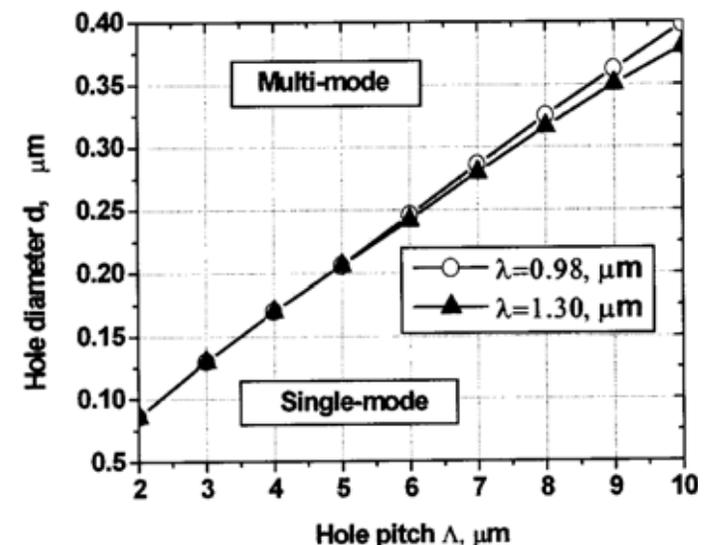
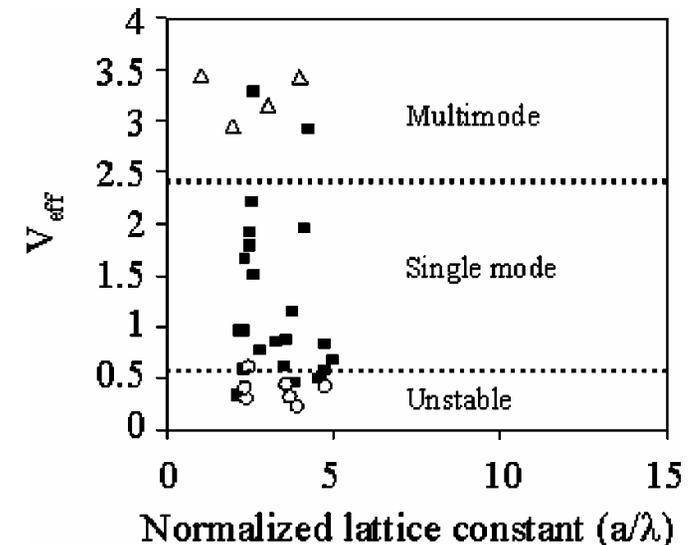
$$B = (\beta_{\text{mode}}^2 - \beta_{\text{core}}^2) / (\beta_{\text{clad}}^2 - \beta_{\text{core}}^2) \quad [3]$$

- β : propagation constant
- $B_{\text{LP}_{01}} < 0.57$ corresponds to the single-mode regime [3]

[1] N. Yokouchi *et al.*, IEEE JSTQE, **9**, p. 1439 (2003)

[2] A. J. Danner *et al.*, APL, **82**, p. 3608 (2003)

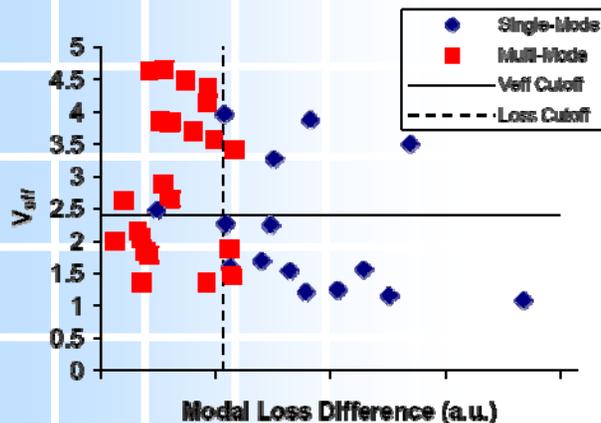
[3] P. S. Ivanov *et al.*, JOSAB, **20**, p. 2442 (2003)



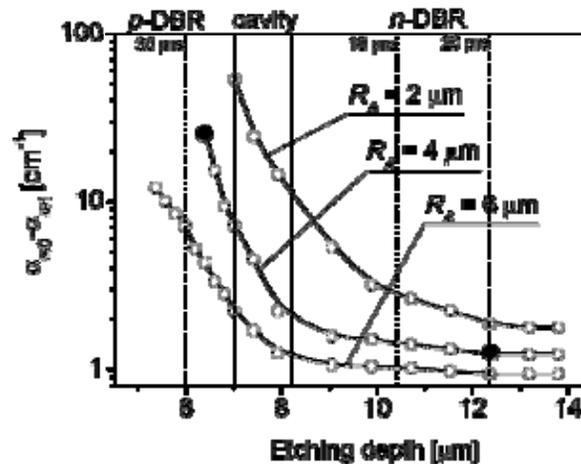
Modal Loss Effect

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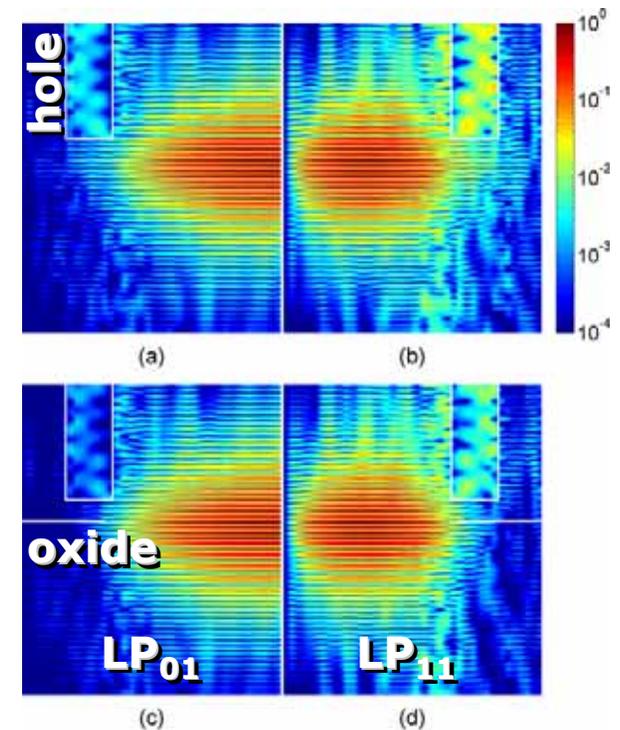
- It was shown experimentally [4] and with simulations that not only the modes' confinement but also their losses influence the single-mode condition.
- Plane wave admittance method [5] and finite element method [6] were applied, both for cold-cavity case.



[4] D. F. Siriani *et al.*, LEOS 2007, WZ3



[5] T. Czystanowski *et al.*, Opt. Express, 15, p. 5604 (2007)



[6] P. Nyakas, JLT, 25, p. 2427 (2007)

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Summary and outlook

Electro-Thermal Models

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Electro-thermal equations [7]

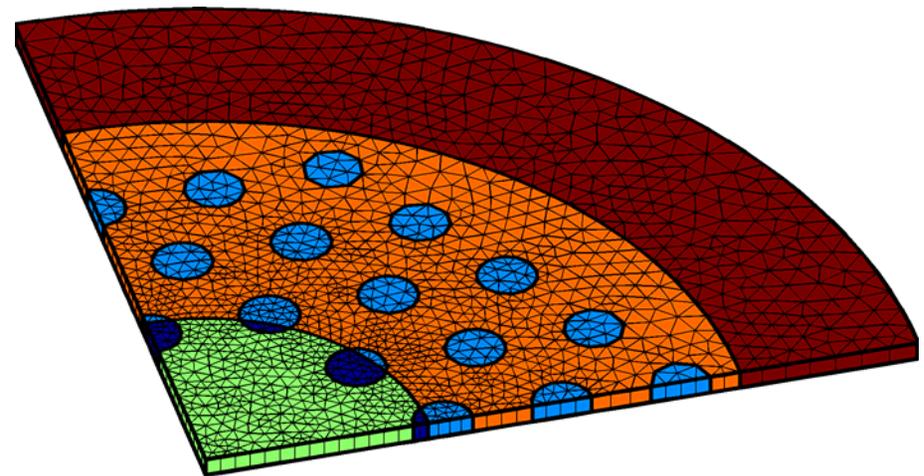
- Laplace equation for the electrostatic potential (Φ)
- anisotropic electric conductivity (σ) for heterojunctions
- stationary solution of the heat conductivity equation with different heat sources (R)

$$\nabla(\sigma \nabla \Phi) = 0 \quad c\rho \frac{\partial T}{\partial t} = \nabla(\kappa \nabla T) + R_{nr} + R_{Joule} + R_{abs}$$

Discretization

- definition of lateral regions by projecting all interfaces to the top plane
- setting up prism elements that respect all interfaces
- integration by finite volume method (box method)

[7] P. Nyakas *et al.*, JOSAB, **23**, p. 1761 (2006)



(green: aperture, orange: implant, deep red: contact and PML, blue: holes)

Material Parameters

Temperature-dependent complex refractivity

- linear dependency of the real part of the index versus temperature
- its coefficient depends on the Al-composition of $\text{Al}_x\text{Ga}_{1-x}\text{As}$, and $\frac{1}{n'} \frac{dn'}{dT}$ varies between $1.25-4 \times 10^{-4}$ 1/K [8]
- we used exponential temperature dependence for the imaginary part that corresponds to free-carrier absorption
- $T_0 = 180$ K was assumed irrespective of the composition due to the lack of more detailed experimental data [9]

Empirical gain function

$$g_i(n, T) = a_0 \ln\left(\frac{n}{n_0}\right) \left\{ 1 - \left[\frac{\lambda_{\text{gain}}(T) - \lambda_i(T)}{\Delta\lambda} \right]^2 \right\}$$

- the gain coefficient (a_0) and the transparency carrier density (n_0) can depend on temperature
- we assumed a parabolic effect of gain-to-cavity detuning

[8] M. Streiff *et al.*, IEEE JSTQE, **9**, p. 879 (2003)

[9] C. H. Henry *et al.*, IEEE JQE, **QE-19**, p. 947 (1983)

3-D Optical Mode Solver

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Scalar and vectorial solutions

- finite volume method was used to solve the scalar Helmholtz equation [7],
- and finite element method was applied to solve the vectorial Helmholtz equation [6]

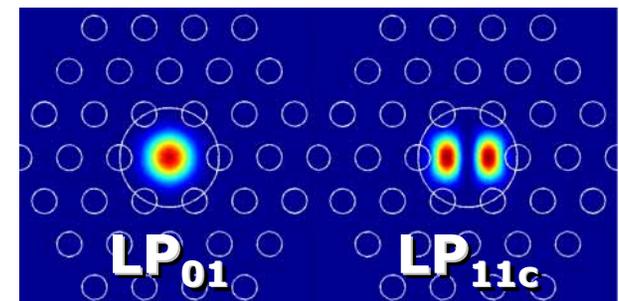
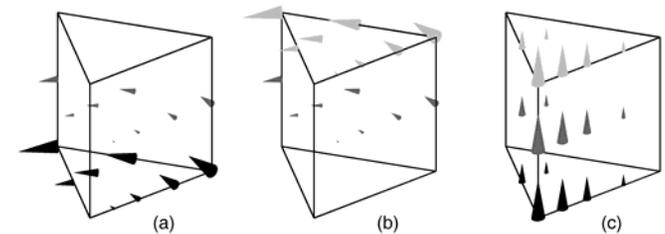
$$F(\mathbf{E}) = \frac{1}{2} \int_V \left[(\nabla \times \mathbf{E}) \mathbf{\Lambda}^{-1} (\nabla \times \mathbf{E}) - \frac{\omega^2}{c_0^2} \mathbf{E} \epsilon \mathbf{E} \right] dV$$

- the scalar solution is used here because of its lower memory and runtime demands

Symmetry boundary conditions

- a 30-degree section is enough to calculate the scalar fundamental and some higher modes of a hexagonal PhC lattice
- a quarter cross-section is needed for LP₁₁ and other higher modes

$$\Delta E = -\frac{\omega^2}{c_0^2} \epsilon E$$



[6] P. Nyakas, JLT, **25**, p. 2427 (2007)

[7] P. Nyakas et al., JOSAB, **23**, p. 1761 (2006)

3-D Optical Mode Solver

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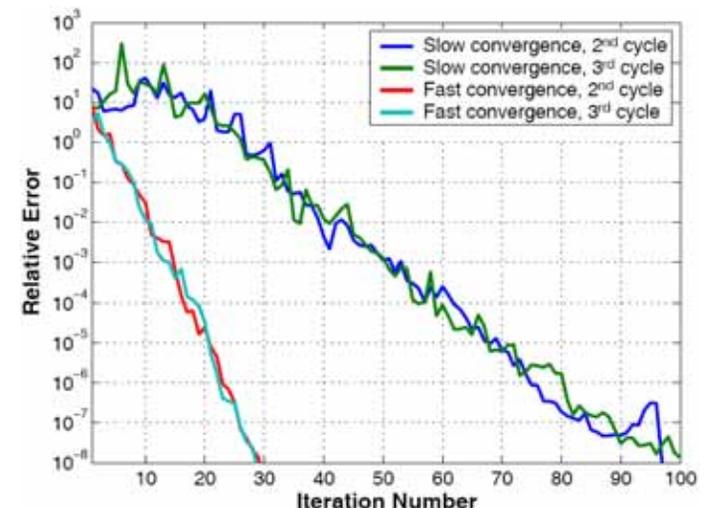
Cold-cavity and active cavity descriptions

- ❑ optical gain is *not* included directly in the complex index when determining the laser resonator modes
- ❑ it is taken into account in the rate equation approach when following the evolution of the modes' power
- ❑ however, local free-carrier absorption loss in the mirrors is included in the refractive index
- ❑ the real part of the index is updated according to the local temperature when calculating the mode profiles

Algebraic problem

- ❑ the complex-symmetric generalized eigenproblem is solved with preconditioned shift-invert iteration
- ❑ the convergence speed and the memory allocation can be tuned with the drop tolerance [7]

[7] P. Nyakas et al., JOSAB, 23, p. 1761 (2006)



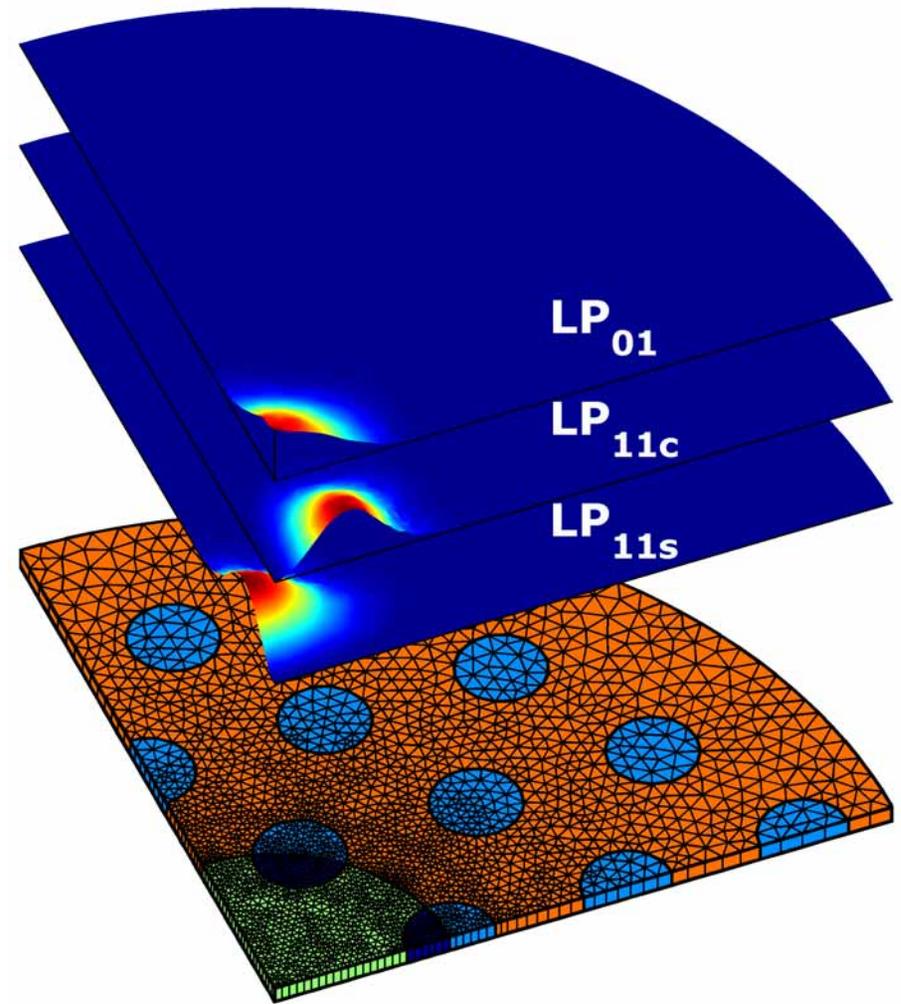
Multi-Mode Rate Equations

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- cover local radiative and nonradiative recombinations, lateral carrier diffusion and spatial hole burning
- the stationary solution is found using an ODE-solver after spatial discretization

$$\frac{\partial n(\mathbf{r},t)}{\partial t} = \frac{\eta j(\mathbf{r},t)}{ed} - An - Bn^2 - Cn^3 + D\Delta n - v_g \sum_i g_i(n) |E_i|^2 S_i$$

$$\frac{dS_i(t)}{dt} = \beta B \int n^2 dV + v_g \left[\int_{QW} g_i(n) |E_i|^2 dV - L_i \right] S_i$$



Coupled Simulations

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Computer demands and simulation times

- ❑ it took about 2-3 GB memory and 20-60 minutes to find the fundamental mode once (30-degree section, 1.5m unknowns)
- ❑ it took about 10-25 GB memory and 10-30 hours to determine LP₁₁ mode once (90-degree section, 4.2m unknowns)

Interpolation technique

- ❑ to accelerate the simulations for varying bias current,
- ❑ we estimated the temperature distribution at discrete points from the electro-thermal equations,
- ❑ and calculated the optical modes under these conditions
- ❑ the transverse mode profiles and the respective quantitative data were then interpolated for interior bias points

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Summary and outlook

Target of Optimization

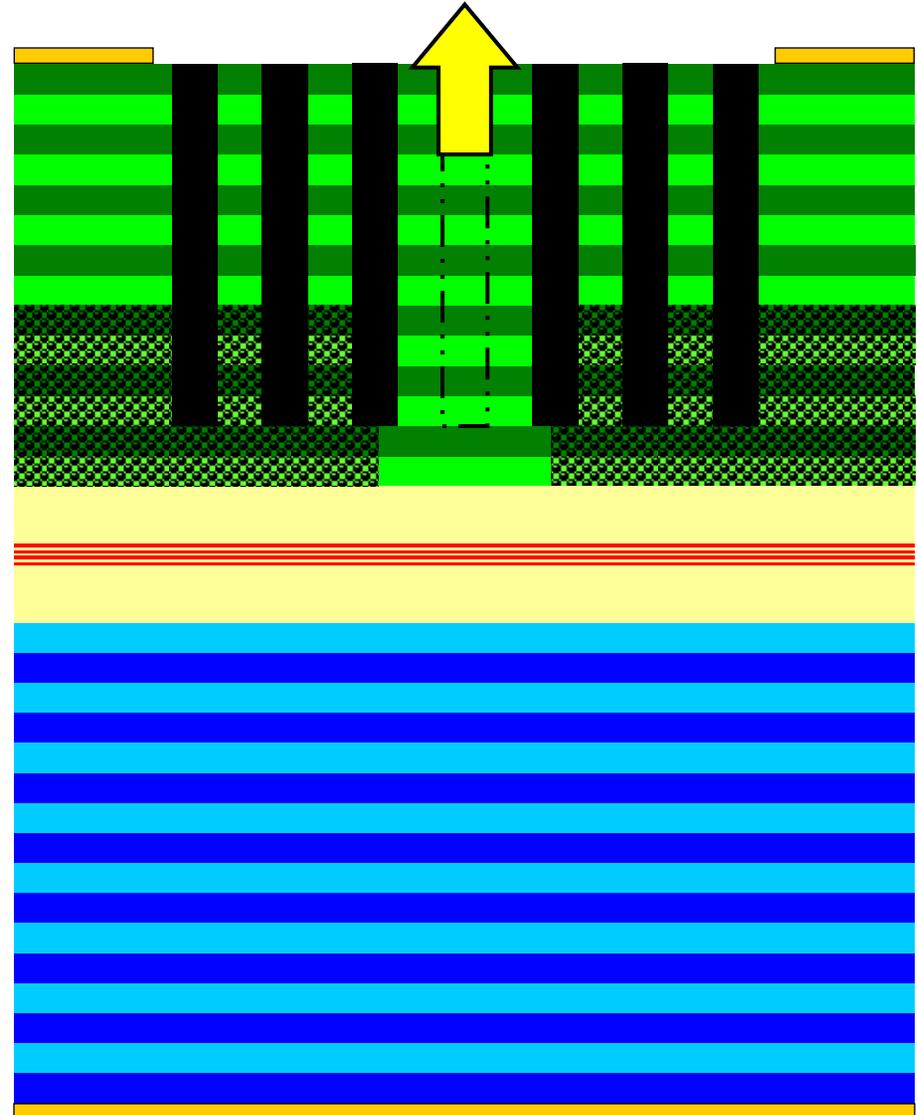
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Fixed parameters

- ❑ design wavelength: 850 nm
- ❑ top-DBR: 30 pairs of p-doped $\text{Al}_{0.9}\text{Ga}_{0.1}\text{As}/\text{Al}_{0.2}\text{Ga}_{0.8}\text{As}$
- ❑ bottom-DBR: 35.5 pairs of n-doped $\text{Al}_{0.9}\text{Ga}_{0.1}\text{As}/\text{Al}_{0.2}\text{Ga}_{0.8}\text{As}$
- ❑ implant thickness: 12 pairs (about $1.5\ \mu\text{m}$)
- ❑ aperture diameter: $7.9\ \mu\text{m}$
- ❑ three rings of holes in a hexagonal pattern around a single defect
- ❑ hole diameter: $0.5\ \Lambda$
- ❑ we consider 2 modes

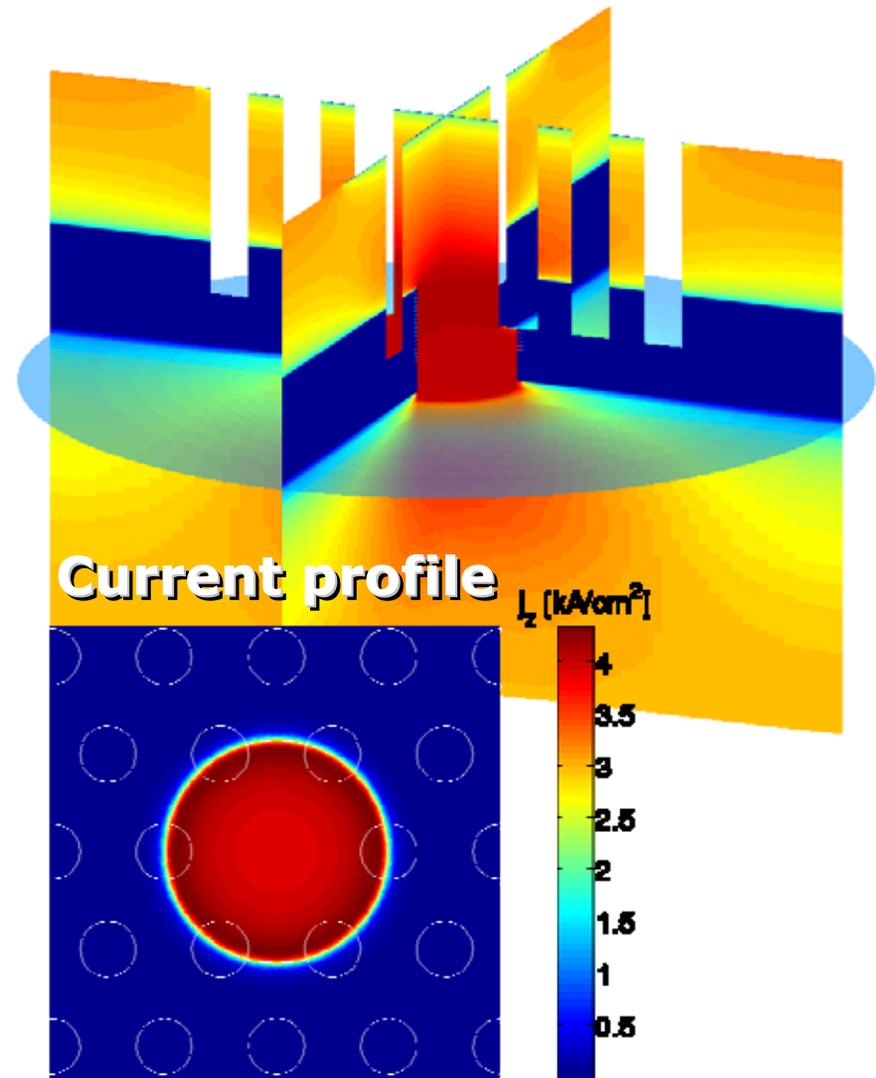
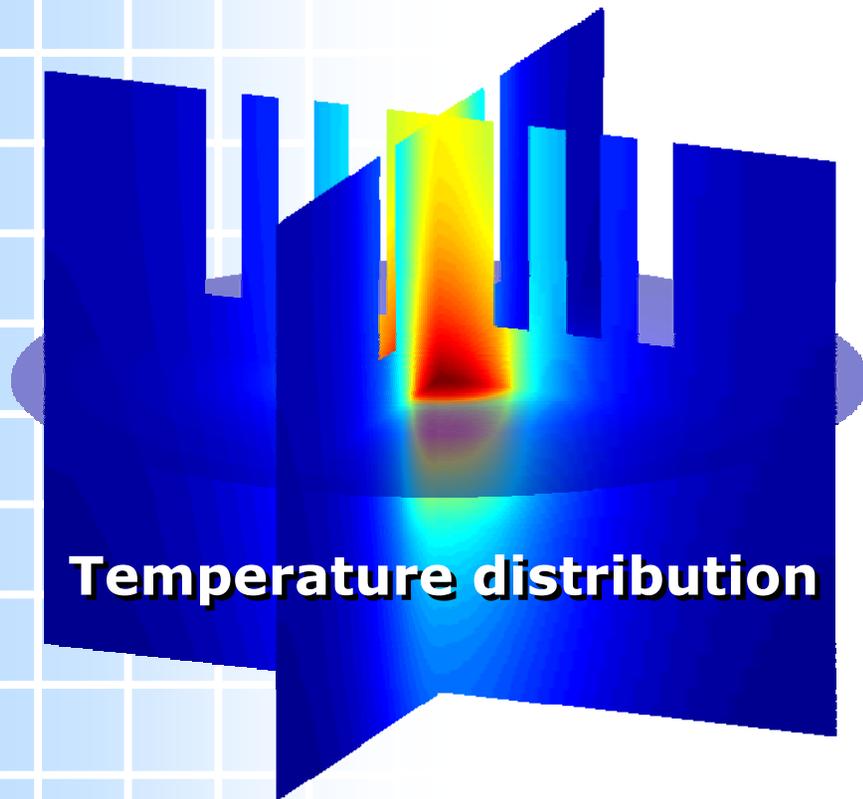
To be optimized

- ❑ lattice constant ($3/4/5\ \mu\text{m}$) to obtain highest single-mode power in fundamental mode



Temperature & Current Profiles

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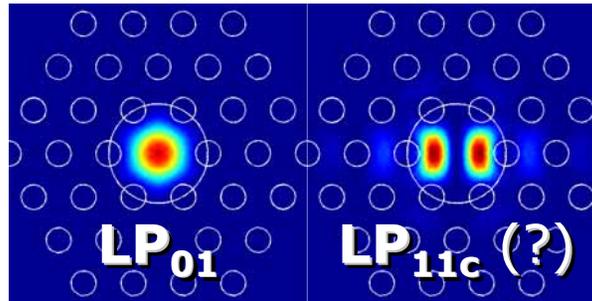


- ❑ The holes do not seem to impact drastically either the heat flow or the current flow.
- ❑ The electric conductivity may degrade also around the holes.

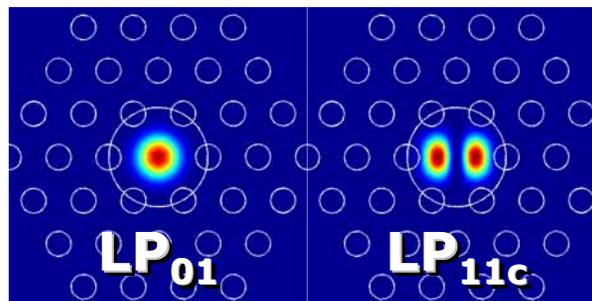
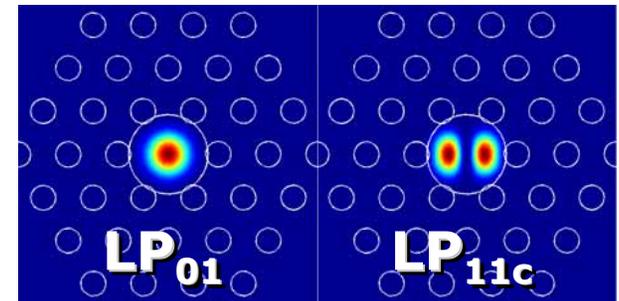
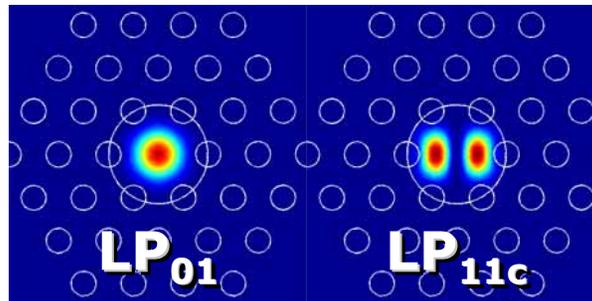
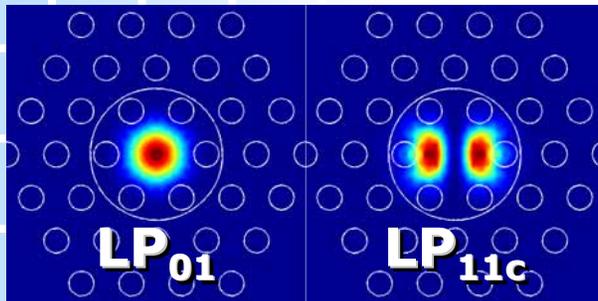
Thermal Lensing

$\Lambda = 4 \mu\text{m}$, 0/10/20 mW heat dissipation

↓
 $\Lambda = 3 \mu\text{m}$,
10 mW heat
dissipation

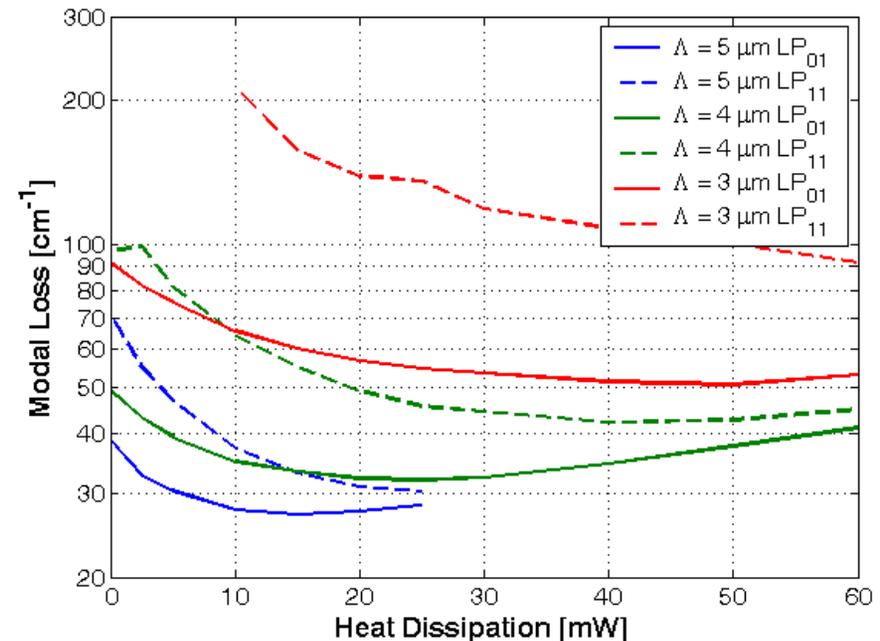
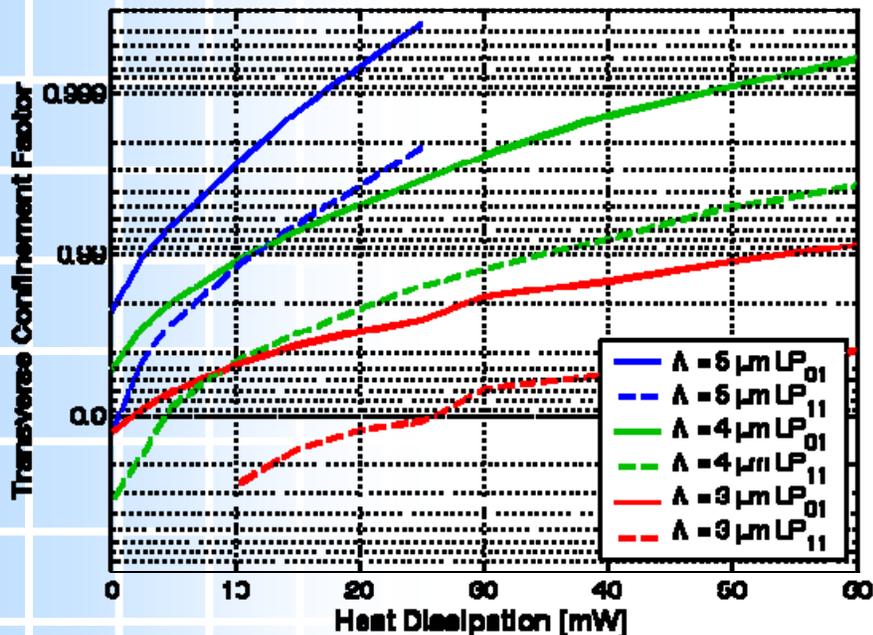


$\Lambda = 5 \mu\text{m}$,
10 mW heat
dissipation ↓



Optical Mode Properties

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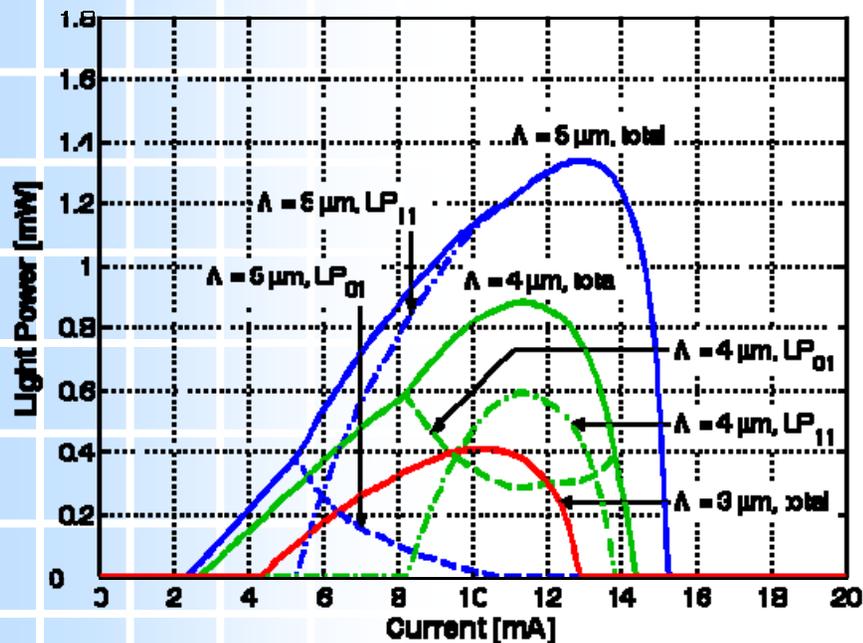


- the transverse confinement increases with larger lattice constant and increasing bias current
- modal discrimination increases as the lattice constant shrinks
- it depends on the structure which of the two effects is dominant

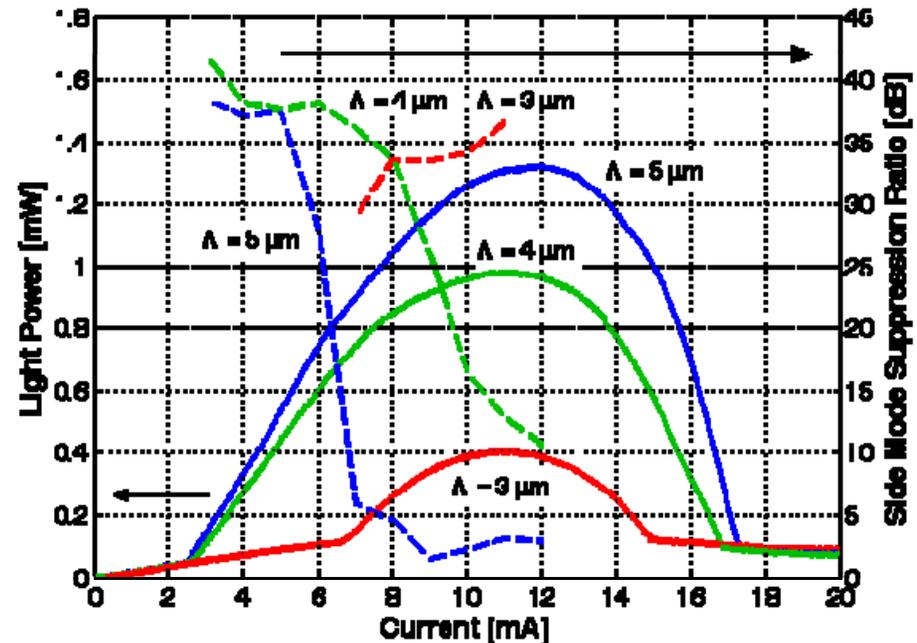
L-I Diagrams

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Calculation



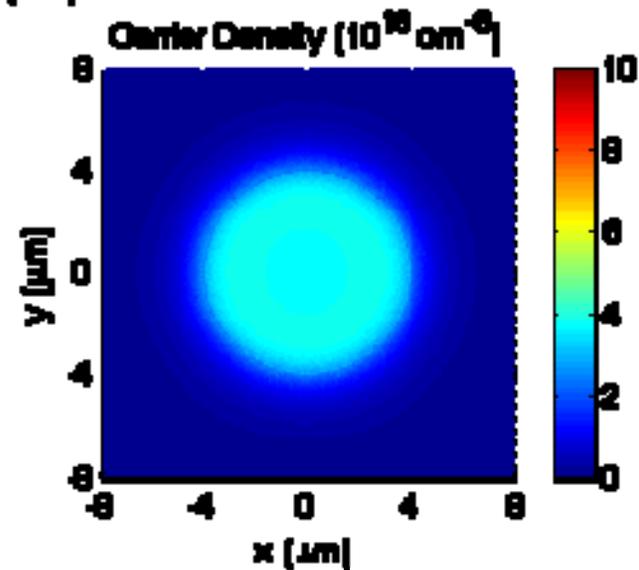
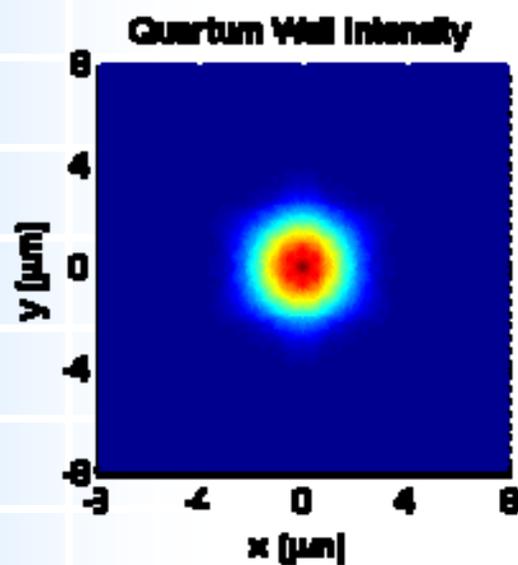
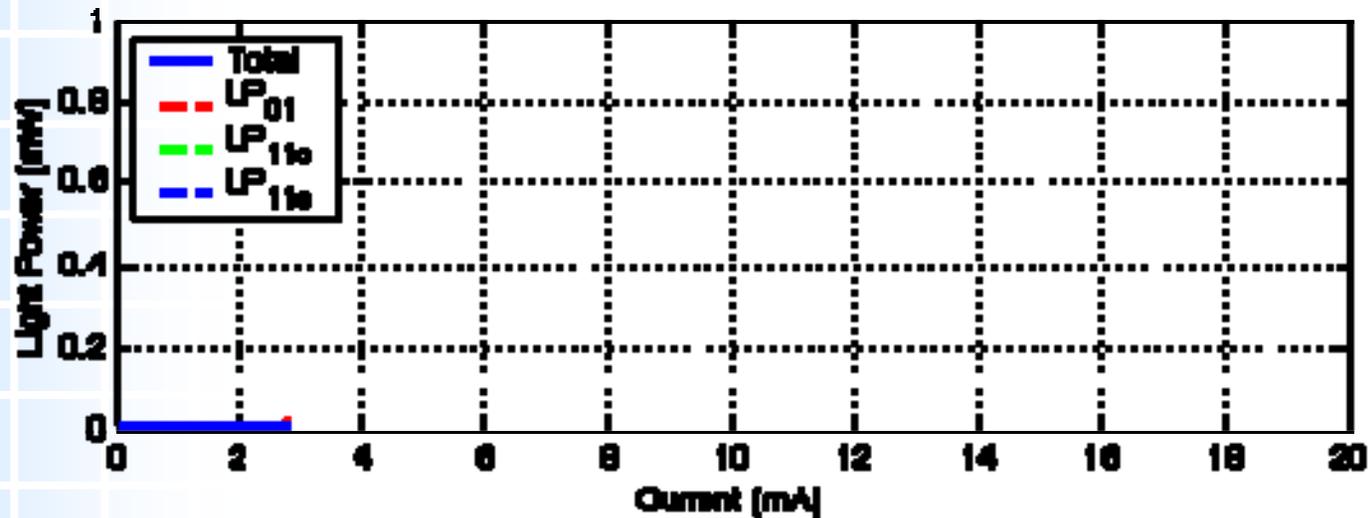
Measurement



- the threshold current increases, and the slope decreases for smaller Λ due to higher optical loss and more wasted current
- single-mode operation can be maintained even if LP₁₁ gets confined, but suffers from high optical loss
- simulation and experiment agree in the optimal $\Lambda = 4 \mu\text{m}$

Spatial Hole Burning

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Summary and Outlook

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Results

- ❑ 3-D coupled simulation model for PhC-VCSELs
- ❑ the confinement and the modal loss discrimination effects have been compared
- ❑ simulation and experiment have shown agreement in the optimal lattice constant that gives highest single-mode power

Further optimization

- ❑ optimal ratio of the diameter of the optical and electrical confinements
- ❑ an oxide aperture provides lower differential resistance,
- ❑ but it can influence the mode confinement and also the losses
- ❑ optimization of small-signal modulation response

Improve the electrical simulation

- ❑ drift-diffusion model aiming realistic semiconductor behavior
- ❑ check the dependence of the differential resistance on the holes' diameters, and determine the effective hole size